

11 **Abstract**

12 This study finds that the correlation between El Niño-Southern Oscillation (ENSO) and
13 the activity of mesoscale oceanic eddies in the South China Sea (SCS) changed around 2004.
14 The mesoscale eddy number determined from satellite altimetry observations using a geometry
15 of velocity vector method was significantly and negatively correlated with the Niño3.4 index
16 before 2004, but the correlation weakened and became insignificant afterward. Further analyses
17 reveal that the ENSO-eddy relation is controlled by two major wind stress forcing mechanisms:
18 one directly related to ENSO and the other indirectly related to ENSO through its subtropical
19 precursor – the Pacific Meridional Modes (PMMs). Both mechanisms induce wind stress curl
20 variations over the SCS that link ENSO to SCS eddy activities. While the direct ENSO
21 mechanism always induces a negative ENSO-eddy correlation through the Walker circulation,
22 the indirect mechanism is dominated by the northern PMM (nPMM) resulting in a negative
23 ENSO-eddy correlation before 2004, but dominated by the southern PMM (sPMM) after 2004
24 resulting in a positive ENSO-eddy correlation. As a result, the direct and indirect mechanisms
25 enhance each other to produce a significant ENSO-eddy relation before 2004, but cancel each
26 other out resulting in a weak ENSO-eddy relation afterward. The relative strengths of the
27 northern and southern PMMs are the key to determining the ENSO-eddy relation and may be
28 related to a phase change of the Interdecadal Pacific Oscillation.

29 **Keywords:** eddy; South China Sea; ENSO; Pacific Meridional Mode; Interdecadal Pacific
30 Oscillation.

31 **1. Introduction**

32 The South China Sea (SCS) is a semi-enclosed sea located in the subtropical western Pacific
33 whose upper ocean circulation is strongly influenced by surface winds in the region (e.g., Fang et
34 al. 1998; Chu et al. 1999; Hu et al. 2000; Fang et al. 2002; Liu et al. 2008; Hu and Wang 2016).
35 The winds produce stress and stress curls that directly drive ocean currents within the basin as
36 well as modulate Kuroshio intrusions into the basin through the Luzon Strait, both of which are
37 the primary factors in determining the circulation pattern in the SCS (e.g., Farris and Wimbush
38 1996; Qu 2000; Qu et al. 2004; Yuan et al. 2006; Wang et al. 2013). The surface wind patterns
39 vary seasonally in association with the seasonal reversal of the East Asian monsoon (e.g., Ding et
40 al. 2004; Wang et al. 2009) and interannually in association with the occurrence of El Niño-
41 Southern Oscillation (ENSO) events (e.g., Chao et al. 1996; B. Wang et al. 2000; Qu et al. 2004,
42 2005; C. Wang et al. 2006; Y. Wang et al. 2006). During ENSO events, anomalous warming and
43 cooling of the sea surface in the central-to-eastern tropical Pacific can disturb the atmospheric
44 circulation resulting in sea surface wind variations over a large part of the Pacific Ocean,
45 including the SCS region (e.g. Zhang et al. 1997; C. Wang et al. 2006; Fang et al. 2006; Wang et
46 al. 2009). The Luzon throughflow can also convey the ENSO influence into the interior SCS
47 (Farris and Wimbush 1996; Qu et al. 2004, 2009; Nan et al. 2015).

48 Surface wind variations over the SCS affect not only the large-scale circulation but also
49 the mesoscale ocean eddies inside the basin. Mesoscale eddies can be observed throughout the
50 SCS (e.g., Soong et al. 1995; Li et al. 1998; Shaw et al. 1999; Wang et al. 2003; Yuan et al. 2007;
51 D. Wang et al. 2008; Xiu et al. 2010; Chen et al. 2011; Lin et al. 2015; Xia and Shen 2015) and
52 have a typical radius of about 100-200 km and a typical lifespan of approximately 8 to 10 weeks
53 (Chen et al. 2011). Despite their chaotic nature, mesoscale ocean eddies play an essential role in

54 transporting water mass, energy, and biochemical substances in the interior SCS and can
55 profoundly impact the regional climate and environment (e.g., Hwang and Chen 2000; Wang et
56 al. 2003; Xiu et al. 2010; Chen et al. 2011). Previous studies have suggested that mesoscale
57 eddies in the SCS can be generated via baroclinic instability of the gyre circulation within the
58 basin (Pedlosky 1982; Wu et al. 1999; L. Wang et al. 2000; Cheng and Qi 2010; Sun et al. 2016)
59 that is forced by the monsoonal flows, the interactions between the monsoonal flows and coastal
60 topography (Wang et al. 2003; G. Wang et al. 2008; Chu et al. 2017), and by the frontal
61 instability (Wang et al. 2003; Gan and Qu 2008) or vorticity advection associated with Kuroshio
62 intrusions (Metzger and Hulburt 2001; Jia and Chassignet 2011; Nof et al. 2011). Each of these
63 eddy generation mechanisms is directly or indirectly related to the prevailing wind stress and
64 wind stress curls over the SCS. Variations in surface winds induced by remote forcings, such as
65 those associated with ENSO, can strengthen or weaken the gyre instability or Kuroshio
66 intrusions to modulate the SCS eddy activity. Mesoscale eddies respond to wind forcing rather
67 rapidly with a response time of about one to several weeks (Chi et al. 1998; G. Wang 2008; Chu
68 et al. 2017).

69 ENSO is a key contributor to the interannual variability in surface winds over the SCS,
70 which in turn should enable ENSO to influence the mesoscale eddy activity in the SCS. However,
71 previous studies were not conclusive concerning the relationship between ENSO and SCS
72 mesoscale eddy activity. Cheng and Qi (2010), for example, found the level of eddy kinetic
73 energy (EKE) in the SCS to be below normal during El Niño events but above normal during La
74 Niña events. They argued that El Niño (La Niña) events can induce anomalous anticyclonic
75 (cyclonic) wind stress curl over the SCS, which weakens (strengthens) the background cyclonic
76 gyre, giving rise to a below (above) normal level of eddy activity during El Niño (La Niña)

77 events. ENSO can also affect the mesoscale eddy activity in the SCS by modulating the Kuroshio
78 intrusions into the basin. Typically, El Niño events weaken the Kuroshio intrusions (Metzger and
79 Hulburt 2001; Metzger 2003; Qu et al. 2004), which reduces frontal instability and/or vorticity
80 advection into the SCS (Cheng and Qi 2010; Nan et al. 2015) and consequently reduces eddy
81 activity in the SCS; and vice versa for La Niña events (Jia and Chassignet 2011; Nof et al. 2011;
82 Nan et al. 2015). ENSO is also suggested to affect the summer monsoon flow which modulates
83 the wind stress curl pattern off Vietnam and thus affects the eddy activity in the SCS (Chu et al.
84 2017). However, there were studies which suggest that the ENSO impact on the SCS eddies is
85 weak or not clear. Xiu et al. (2010) employed a numerical model to examine the relationship
86 between ENSO and the eddy number and found no direct correlation. Chen et al. (2011) found
87 no obvious correlation between ENSO and the SCS eddy number they identified from satellite
88 observations. The different findings on the ENSO-eddy relationship may be caused by the
89 different datasets or identification methods used in the studies or may be due to the existence of
90 decadal changes in the ENSO-eddy relationship.

91 Due to their random and chaotic characteristics, mesoscale eddies are not easy to directly
92 identify from observations. In this study, we employ an automatic eddy identification method
93 based on the geometry of velocity (GV) field (Nencioli et al. 2010) to determine the mesoscale
94 eddy number in the SCS from satellite altimetry observations. The GV method has been shown
95 to determine the eddy number with reasonably good accuracy (Lin et al. 2015; Xia and Shen
96 2016), compared with other methods (Wang et al. 2003; Xiu et al. 2010; Chen et al. 2011). We
97 then use the identified eddy number to examine the relationship between ENSO and SCS
98 mesoscale eddy activity. We discovered that the ENSO-eddy relationship is not stationary but
99 changes from decade to decade. Furthermore, we find the mesoscale eddy activity in the SCS is

100 also affected by the remote forcing from an ENSO precursor – the Pacific Meridional Modes
101 (PMMs) in the subtropical Pacific of both hemispheres. The subtropical Pacific influences,
102 which were not emphasized by previous studies, are a reason why the ENSO-eddy relationship
103 changes over the decades. The large-scale dynamical mechanisms that link SCS mesoscale
104 eddies to ENSO and the PMMs are then identified and explained.

105 This paper is organized as follows: Data and the method for identifying mesoscale eddies
106 used in this paper are introduced in Section 2. The wind stress curl modes and associated large-
107 scale atmospheric and oceanic anomalies are described in Sections 3-4. The possible
108 mechanisms for the decadal variation in ENSO-eddy relationship are discussed and analyzed in
109 Section 5. Section 6 summarizes the findings and implications from this study.

110 **2. Data and Methods**

111 **a. Datasets**

112 In this study, we used the sea level anomaly (SLA) and geostrophic current anomaly data
113 from the Archiving, Validation, and Interpretation of Satellite Oceanographic (AVISO) mission
114 to identify eddies in the SCS. The merged daily AVISO data are available from January 1993 to
115 October 2014 and have a resolution of $1/4^\circ$ longitude by $1/4^\circ$ latitude. The altimeter data contain
116 aliases due to the shelf, tidal, and internal waves over the shelf area (Yuan et al. 2006), which
117 can introduce errors in eddy detection algorithms. Therefore, the SLA data in waters with depths
118 shallower than 200 m are not used in this study. Also used are the monthly sea surface wind, sea
119 surface temperature (SST), and sea level pressure (SLP) data during the period 1993-2014. The
120 sea surface wind data are provided by Cross-Calibrated Multi-Platform (CCMP) project that
121 includes cross-calibrated satellite winds derived from a combination of many radar

122 scatterometers (Atlas et al. 2011). The monthly CCMP wind product also has a resolution of $1/4^\circ$
123 longitude by $1/4^\circ$ latitude. The monthly SST and SLP data are from the ERA-Interim Reanalysis
124 of European Centre for Medium-Range Weather Forecasts, which has a resolution of $1/4^\circ$
125 longitude by $1/4^\circ$ latitude grid (Dee et al. 2011). To calculate wind stress curl, we first used the
126 bulk formulation of Trenberth et al. (1990) to calculate zonal and meridional components of
127 wind stress τ_x and τ_y , respectively. We then applied Stokes' theorem to obtain the vertical
128 component of $\text{curl}_z(\tau) = \nabla \times \tau$, which is the surface wind stress curl. To represent ENSO activity,
129 we obtained the monthly Niño3.4 index from the National Oceanic and Atmospheric
130 Administration (NOAA). We also used the Interdecadal Pacific Oscillation (IPO) index from
131 NOAA, and calculated the North Pacific Oscillation (NPO) index based on the previous study by
132 Yu et al. (2012). Indices to represent the strengths of the northern and southern PMMs were also
133 used in this study. Following Zhang et al. (2014), we calculated the monthly northern and
134 southern PMM indices as the SST anomaly average over the northeast Pacific ($21\text{-}25^\circ\text{N}$, 138-
135 142°W) and the southeast Pacific ($19\text{-}15^\circ\text{S}$, $103\text{-}107^\circ\text{W}$), respectively. Anomalies in this study
136 are defined as the deviations from the climatological seasonal cycle for the period 1993-2014
137 after removing the linear trend.

138 **b. Identification of mesoscale oceanic eddies**

139 The GV method examines the geometry of velocity vectors to identify eddies (Nencioli et
140 al. 2010). A mesoscale eddy is defined as a flow feature containing a consistent sense of rotation
141 relative to a center of minimum speed and is surrounded by an enclosed streamline (as shown in
142 Fig. 1). When applying the GV method to detect an eddy, there are two major steps in the
143 calculation: one is to detect the eddy center, and the other one is to determine the enclosed
144 streamline that corresponds to the eddy center.

145 In searching for a potential eddy center, we need to localize the search area. Two
146 parameters a and b are used to localize the search area and set the minimum size of eddy
147 detectable. Here the parameter a defines how many points away from a chosen center point will
148 be examined, and the parameter b defines the horizontal size (in grid points) of the area used to
149 calculate the local minimum velocity. Their values have to be tuned based on the data resolution
150 to optimize the performance of the GV method (Nencioli et al. 2010). The performance can be
151 assessed by the “success detection rate” ($SDR = \frac{N_c}{N_i}$) and the “excess detection rate” ($EDR = \frac{N_e}{N_i}$),
152 where N_c is the number of eddies detected by both the expert and the detection method, N_i is the
153 number of eddies identified by the expert only, and N_e is the number of eddies detected by the
154 eddy detection method only. Due to the limited horizontal resolution of the altimeter data, only
155 eddies with a radius larger than 25 km ($1/4^\circ$ grid) can be resolved. Earlier studies have found that
156 linearly interpolating the velocity field to higher-resolution grids can significantly improve the
157 SDR and reduce the EDR (Liu et al. 2012; Lin et al. 2015; Xia and Shen 2015). To achieve the
158 minimum size of eddy detectable standard while at the same time conserving computation
159 resources, we choose to linearly interpolate the velocity field to a finer sub-grid of $1/20^\circ \times 1/20^\circ$.
160 This interpolation enables the minimum detectable radius to be reduced to 25 km. The parameter
161 values used here ($a=5$ and $b=4$) are adapted from Xia and Shen (2015) who showed that these
162 values produced the highest SDR and the lowest EDR when applying the GV method to the same
163 AVISO dataset used in this study.

164 The outer edge of the eddy determined by the streamlines is calculated from the stream
165 function ψ : $u' = -\frac{\partial\psi}{\partial y}$ and $v' = \frac{\partial\psi}{\partial x}$. Here, u' and v' are the zonal and meridional components of
166 the sea-surface geostrophic current anomaly, respectively. The edge of an eddy is defined as the

167 outmost enclosed streamline around the eddy center provided that the velocity magnitudes from
168 the enclosed streamline along an arbitrary radius to the eddy center decrease.

169 This GV method enables us to count the daily number of mesoscale eddies in the SCS.
170 Figure 1 shows, as an example, the locations of anticyclonic eddies (AEs) and cyclonic eddies
171 (CEs) identified by the GV method for January 26, 1993. On this particular day, there are a total
172 of 15 mesoscale eddies in the SCS, which includes 9 CEs and 6 AEs. The eddies on this day are
173 distributed more or less evenly across the SCS basin, which is in general consistent with mean
174 climatology of eddy probability reported in earlier studies (e.g., Chen et al. 2011). It should be
175 noted here that only SLA differences larger than 2 cm can be observed accurately by the satellite
176 (Pujol et al. 2016). In order to avoid errors introduced by the observations, we omitted any eddy
177 with an amplitude smaller than 2 cm from our results (e.g., the two AEs off Hainan in Fig. 1).

178 **3. Relationships between ENSO and SCS mesoscale eddies**

179 We examine the eddy-ENSO relationship first by showing in Fig. 2a the time series of the
180 monthly values of the eddy number in the SCS and the Niño3.4 SST index. The monthly values
181 of the eddy number were calculated as the sum of the daily eddy number identified by the GV
182 method and ranged between 400 and 500 eddies (per month) during 1993-2014. The immediate
183 impression from the figure is that there is a tendency for the eddy number to be out-of-phase
184 with the Niño3.4 index. This indicates that the eddy number in the SCS decreases during El Niño
185 events but increases during La Niña events. This negative ENSO-eddy relationship is consistent
186 with the findings of some previous studies (e.g., Cheng and Qi 2010; Chen et al. 2011; He et al.
187 2016; Chu et al. 2017). We have also used three other identification methods, i.e., the sea level
188 anomaly (SLA) based method (Wang et al., 2003), the Okubo-Weiss (OW) approach method

189 (Okubo, 1970; Weiss, 1991) and the Hybrid (HY) method (Yi et al., 2014), to calculate the
190 mesoscale eddy number in the SCS, and then to recalculate the ENSO-eddy relationship; the
191 results obtained (not shown) are similar to those in Fig. 2b.

192 However, the correlation does not seem to be stationary throughout the entire analysis
193 period. In particular, the out-of-phase relationship becomes less evident after 2004. This initial
194 analysis indicates that there may be a change in the ENSO-eddy relationship around 2004. To
195 confirm this and to more precisely determine the timing of the change, we show in Fig. 2b a
196 sliding correlation between the eddy number and the Niño3.4 index using a 10-year moving
197 window. This figure confirms that the ENSO-eddy relationship changed around 2004 from being
198 significantly and negatively correlated before to non-correlated afterward. The eddy number
199 shows a stronger negative correlation with the Niño3.4 index that passes the 95% significance
200 level only during the pre-2004 era. Therefore, in the rest of this study, we focus on understanding
201 how the ENSO-eddy relationships are established before and after 2004.

202 Surface wind stress curl is one of the most relevant generation mechanisms for mesoscale
203 oceanic eddies and needs to be analyzed in order to understand the ENSO-eddy relationships. To
204 identify the leading modes of variation in surface wind stress curl over the SCS, we perform an
205 empirical orthogonal function (EOF) analysis to the curl anomalies over the SCS (5°N-25°N and
206 105°E-123°E). The two leading EOF modes (EOF1 and EOF2, hereafter) explain, respectively,
207 21.3% and 12.8% of the total variance and distinguish themselves from the rest of the EOF
208 modes (not shown). Therefore, we focus only on these two modes.

209 The spatial pattern of EOF1 (Fig. 3a) consists of a basin-scale monopole of wind stress curl
210 anomalies with an elongated band extending from the center (115°E, 15°N) of the analysis box
211 towards a region to the northwest of Luzon Island. In its positive phase, EOF1 is characterized

212 by an anomalous anticyclone covering almost the entire SCS basin. This wind stress curl pattern
213 can spin down the gyre-scale circulation in the SCS, which is cyclonic over the entire SCS
214 during boreal winter (Fang et al. 2006; Y. Wang et al. 2006; Cheng and Qi 2010) and cyclonic
215 north of 12°N during boreal summer (Fang et al., 2006; G. Wang et al. 2008). Therefore, the
216 EOF1 pattern (in its positive phase) can weaken the SCS gyre-scale circulation and its associated
217 baroclinic instability resulting in a reduction in the mesoscale eddy activity in the SCS. As for
218 EOF2 (Fig. 3b), its spatial pattern exhibits a meridional dipole of wind stress curl anomalies with
219 a slight northwest-to-southeast tilt. The northwest lobe of the dipole is located to the east of
220 Hainan Island, while the southeast lobe is located to the west of Luzon Island. Associated with
221 this anomalous curl pattern, an anticyclonic wind stress anomaly pattern occupies most of the
222 SCS. The negative wind stress curl anomalies are stronger and occupy most of the SCS,
223 especially north of 12°N, while the positive wind stress curl anomalies are weaker and occupy
224 the regions to the south. This EOF pattern (in its positive phase) can spin down the summer gyre-
225 scale circulation as well as a large part of the winter gyre-scale circulation resulting in a
226 weakening of SCS mesoscale eddy activity during both seasons. These two EOF patterns can
227 also affect the SCS eddy activity by modulating the Kuroshio intrusion, which is another major
228 generation mechanism for the mesoscale eddies in the SCS (Metzger and Hurlburt 2001; Jia and
229 Chassignet, 2011; Nof et al. 2011). In their positive phase, both EOF1 and EOF2 have positive
230 wind stress curl anomalies southwest of Taiwan that lower local sea surface heights through
231 Ekman transport, and negative wind stress curl anomalies northwest of Luzon that elevate local
232 sea surface heights (Metzger and Hurlburt, 2001; Liang et al., 2008; Hsin et al., 2012; Wu and
233 Hsin, 2012). The resulting meridional pressure gradient across the Luzon Strait can weaken the
234 Kuroshio intrusion into the SCS (Qu, 2000; Metzger and Hurlburt, 2001; Liang et al., 2008; Hsin

235 et al., 2012; Wu and Hsin, 2012), reducing mesoscale eddy activity in the SCS. Therefore, both
236 the EOF1 and EOF2 modes tend to reduce (increase) eddy production over a large part of the
237 SCS by influencing the SCS gyre-scale circulation and Kuroshio intrusions during their positive
238 (negative) phase.

239 We next examine how these two leading modes are related to ENSO. Figure 4 shows the
240 principal components of these two EOFs (i.e., PC1 and PC2) and their 10-year sliding correlation
241 with the Niño3.4 index. As shown in Fig. 4b, it is interesting to find that these two modes show
242 similar positive correlations with Niño3.4 index before 2004 but dramatically opposite
243 correlations with Niño3.4 index after 2004. The EOF1 mode of wind stress curl anomalies
244 maintains its positive correlation with the Niño3.4 index throughout the analysis period, but the
245 EOF2 mode changes from being positively correlated with Niño3.4 index before 2004 to
246 negatively correlated afterward. Therefore, these two EOF modes respond similarly to ENSO
247 before 2004 to enable the El Niño (La Niña) to weaken (strengthen) the production of SCS
248 mesoscale eddies, but respond oppositely after 2004 to cancel each other out resulting in a weak
249 relationship between SCS eddies and ENSO. The change in the correlation seen in Fig. 2b
250 between ENSO and the SCS eddy number around 2004 can be a result of the enhancement and
251 cancelation between these two EOF modes before and after that time, respectively.

252 **4. Large-scale atmospheric and oceanic anomalies associated with modes of SCS wind** 253 **stress curl variability**

254 We then examine the large-scale atmospheric circulation and SST anomalies associated
255 with these two leading modes by regressing SST, SLP, and surface wind anomalies onto PC1
256 and PC2 during the period before 2004 (1993-2003; P1) and the period afterward (2004-2014;

257 P2). The SST regressions onto PC1 during P1 (Figs. 5a-d) and P2 (Figs. 5e-h) are both
258 dominated by a typical evolution of El Niño (Rasmusson and Carpenter 1982). The SST
259 anomalies originate first off the South American Coast and spread westward along the equatorial
260 Pacific. The regressions of SLP onto PC1 (Figs. 6a-d and 6e-h) also reveal several typical
261 features associated with a developing ENSO event. One of them is the Southern Oscillation
262 pattern over the tropical Pacific that is characterized by negative SLP anomalies over the tropical
263 eastern Pacific and positive SLP anomalies over the tropical western Pacific. When we zoom in
264 the Lag 0 regressions around the SCS region (Figs. 7a-b), we can clearly see that the positive
265 SLP anomalies over the SCS are part of the western Pacific center of the Southern Oscillation
266 during both P1 and P2 (cf. Figs. 7a, b to 6c, g). The positive SLP anomalies over the SCS have a
267 local center over the Philippine Sea, which induces an anticyclonic wind stress pattern over the
268 entire SCS that resembles the EOF1 mode (cf. Figs. 7a-b to 3a). These regression analyses
269 indicate that the EOF1 mode of wind stress curl variations over the SCS is part of the Southern
270 Oscillation that accompanies the developing ENSO during both P1 and P2. This explains why
271 the EOF1 mode maintains a stationary positive correlation with the Niño3.4 index throughout the
272 analysis period.

273 As mentioned above, the EOF2 mode exhibits a remarkable change in its relationship with
274 the Niño3.4 index around 2004. During the pre-2004 period (i.e., P1), the SST regression (Figs.
275 5i-l) is dominated by a northern Pacific Meridional Mode (nPMM) (Chiang and Vimont 2004)
276 like pattern that features warm SST anomalies spreading from the North American Coast to the
277 tropical central Pacific and cool SST anomalies that persist in the tropical eastern Pacific. After
278 PC2 and the nPMM reaches their peak phases at Lag 0, the warm SST anomalies at the tropical
279 central Pacific continue to develop into an El Niño event. This El Niño resembles more closely

280 the Central Pacific El Niño (Yu and Kao 2007; Kao and Yu 2009) than the conventional Eastern
281 Pacific El Niño. This is consistent with the suggestion that the nPMM is a precursor of ENSO
282 (e.g., Chang et al. 2007; Yu et al. 2010; Yu and Kim 2011; Yu et al. 2017). As such, PC2 leads
283 the development of the El Niño, in contrast to PC1 that develops together with the El Niño.
284 Therefore, the EOF2 mode of the surface wind stress curl variations in the SCS is associated
285 with a precursor of ENSO before 2004 (i.e., during P1). The positive correlation between PC2
286 and the Niño3.4 index in Fig. 4c during this sub-period does not represent an SCS response to
287 the El Niño but rather a joint connection of the EOF2 and ENSO with the nPMM. We show in
288 Fig. 8 the lead-lag correlations between the two PCs and the Niño3.4, nPMM, and sPMM indices
289 for the two periods (P1 and P2). The figure shows that PC2 has a larger simultaneous (Lag 0)
290 correlation coefficient with the nPMM index (0.37) than with the Niño3.4 index (0.08) during
291 the P1 sub-period (Fig. 8c). In contrast, PC1 has a larger simultaneous correlation with the
292 Niño3.4 index (0.37) during this period. Our analyses indicate that the negative correlation
293 between the SCS eddy activity and ENSO can be established directly through ENSO (i.e., the
294 EOF1 mode) and indirectly through an ENSO precursor – the northern PMM (i.e., the EOF2
295 mode) before 2004.

296 The SLP regression onto PC2 during the pre-2004 period (Figs. 6i-l) is dominated by an
297 NPO pattern (Walker and Bliss 1932; Rogers 1981; Linkin and Nigam 2008), which is
298 characterized by an out-of-phase variation between the Aleutian Low and the Pacific Subtropical
299 High. The correlation coefficient between PC2 and the NPO index is 0.75 during P1. Recent
300 studies have suggested that the NPO can induce SST anomalies off Baja California via
301 anomalous surface heat fluxes, which then spread southwestward via subtropical Pacific
302 atmosphere-ocean coupling into the equatorial Pacific to trigger the development of ENSO in the

303 tropical central Pacific (Kao and Yu 2009; Yu et al. 2010; Yu and Kim 2013; Yu et al. 2017).
304 This sequence of events is similar to those associated with the positive nPMM pattern and the
305 ENSO onset shown in Figs. 5i-l. The regressed surface wind anomalies (particularly at Lag 0;
306 Figs 5k and 6k) show that an anomalous surface cyclone forms over the Western North Pacific
307 (covering the region from Japan to Taiwan) in association with the nPMM. This cyclonic
308 anomaly is a Gill-type response to the positive SST anomalies associated with the nPMM that
309 has been mentioned in previous studies of the PMM (e.g., Wang et al. 2012; Zhang et al. 2016).
310 This anomalous cyclonic circulation induces anomalous northerly winds over the SCS similar to
311 those in the EOF2 pattern (cf. Fig. 7c to 3b).

312 During the post-2004 period (i.e., P2), the SST regression onto PC2 (Figs. 5m-p) is very
313 different from that obtained during the pre-2004 period. The regressed anomalies are dominated
314 by a southern Pacific Meridional Mode (sPMM) (Zhang et al. 2014) in the Southeastern Pacific.
315 The sPMM is the southern-hemispheric analog of the nPMM and is characterized by SST
316 anomalies extending from the Peruvian Coast toward the equatorial central Pacific. As shown in
317 Fig. 8d, PC2 has a larger simultaneous correlation coefficient with the sPMM index (-0.22) than
318 with the nPMM index (-0.07) during the P2 sub-period. The sPMM is capable of developing into
319 the deep tropics through its connection with cold tongue ocean dynamics (e.g., mean advection)
320 (Zhang et al. 2014). As shown in Figs. 5m-p, cold SST anomalies developed at the equator
321 together with the evolution of the negative phase of the sPMM. As a result, PC2 also shows a
322 large negative correlation with the Niño 3.4 index during this period. It should be noted that the
323 positive values of the PC2 are associated with a “cold” phase of the sPMM and a La Niña during
324 P2 but associated with a “warm” phase of the nPMM and an El Niño during P1. This explains
325 why the correlation between PC2 and the Niño3.4 index changed from positive before 2004 to

326 negative afterward. The SLP anomalies regressed onto PC2 during P2 (Figs. 6m-p) are
327 dominated by positive values over the Southeastern Pacific, a typical pattern associated with the
328 cold sPMM pattern (cf. Fig. 1f of Zhang et al. 2014). The regressed surface wind anomalies
329 indicate an anomalous anticyclonic circulation formed between 130-180°E as a Gill-type
330 response to the cold SST anomalies associated with the sPMM (Figs. 5g and 6g). As part of a
331 Rossby wave response, an anomalous cyclonic circulation forms further to the west over the
332 Western North Pacific. The anomalous cyclone extends from Japan into the Philippine Sea and
333 induces an anomalous northerly pattern in the SCS that is similar to the EOF2 mode. This
334 similarity can be better seen by focusing on the anomalies around the SCS (Fig. 7d). It is very
335 interesting to find that a similar wind stress anomaly pattern in the SCS can be produced as a
336 direct Gill-type response to the nPMM during P1 or as part of a Rossby wave train associated
337 with the Gill-type response to the sPMM during P2.

338 To further confirm that EOF2 is indeed related to the nPMM before 2004 (i.e., P1) but to
339 the sPMM afterward (i.e., P2), we show in Fig. 9 the wind and SLP anomaly patterns regressed
340 onto the nPMM and sPMM indices during the two sub-periods. The figure shows that, during P1,
341 the nPMM-regressed wind and SLP anomaly patterns are similar to those regressed onto the PC2
342 during P1 (cf. Figs. 6k and 9a), with a pattern correlation coefficient of 0.76. In both sets of
343 regressions, surface northerlies prevail south of 12°N throughout the SCS. However, the nPMM
344 regressions during P2 are very different from the PC2 regressions during that period (cf. Figs. 6g
345 and 9b), with a pattern correlation coefficient between of only -0.23. Besides having a very
346 different SLP anomaly pattern from the PC2 regressions, the nPMM regressions show very weak
347 surface wind anomalies over the SCS during this period. In contrast, sPMM-regressed wind and
348 SLP anomaly are similar to those regressed onto PC2 during P2 (cf. Figs. 6g and 9d) but not

349 during P1 (cf. Figs. 6k and 9c). Their pattern correlations reach 0.42 during P2 but only 0.22
350 during P1. Since the positive phase of EOF2 is related to the negative phase of the sPMM, the
351 signs of the regression patterns in Figs. 9c and 9d have been reversed to aid the comparison with
352 Figs. 6g and 6k. The sPMM regressions show prevailing northerlies throughout the entire SCS
353 during P2 but only to the north of 12°N during P1.

354 Thus, our analyses suggest that the EOF1 mode of wind stress curl variations in the SCS is
355 directly forced by ENSO via the Southern Oscillation throughout the analysis period, whereas
356 the EOF2 pattern is forced by the nPMM before 2004 and by the sPMM after 2004. Also, the
357 nPMM links the positive phase of the EOF2 to El Niño, while the sPMM links the positive EOF2
358 to La Niña. As a result, when an El Niño event occurred during the pre-2004 period, it was
359 associated with a positive phase of EOF1 and a positive phase of the EOF2 to together produce
360 anticyclonic wind stress curl anomalies over the SCS. This weakened wind stress curl and thus
361 curbed the production of mesoscale eddies in the SCS. In contrast, when an El Niño event
362 occurred during the post-2004 period, it was associated with a positive phase of the EOF1 and a
363 negative phase of the EOF2, and the wind stress curl anomalies over the SCS tended to be weak
364 as a result and so there was a weak ENSO impact on SCS eddy activity during this sub-period.
365 The differing relationship between the SCS eddy number and the ENSO before and after 2004 is
366 caused by the differing associations of the EOF2 mode with the northern and southern PMMs
367 around that time. To further confirm the changed relationships between the SCS eddy number
368 and the two PMMs, we show in Fig. 10 the 10-year sliding correlations between monthly eddy
369 number anomaly and the nPMM and sPMM indices. This figure clearly shows that the SCS eddy
370 number was more strongly correlated with the nPMM during the period before 2004 but more
371 correlated with the sPMM afterward. This result adds support to our finding on the relative

372 importance of the northern and southern PMMs for the SCS eddy activity changes from decade
373 to decade.

374 **5. Influences of the IPO on the recent ENSO-eddy relation changes**

375 A key question to ask is what causes the EOF2-PMM relation to change around 2004? We
376 notice that 2004 is close to the time when the IPO switches from its positive to negative phase
377 (Fig. 11). The IPO's phase change may be the reason for the change in EOF2-PMM relations
378 around 2004. The SST and SLP anomalies regressed onto the IPO index during the sub-periods
379 before and after 2004 are shown in Fig. 12. It shows that the ENSO-like SST variability
380 associated with the IPO (Figs. 12a and 12c) is associated with large SLP anomalies over both the
381 Northeastern and Southeastern Pacific (Figs. 12b and 12d) during both periods. These are the
382 regions where the nPMM and sPMM are located. The figure indicates that SLP anomalies over
383 the sPMM region increased in magnitude from P1 to P2. Figure 12 suggests that the SLP
384 anomalies produced by the IPO change the strengths of the background trade winds in these two
385 regions resulting in a difference in the relative importance of the nPMM and sPMM before and
386 after 2004. This possible modulation effect of the IPO on the nPMM and sPMM requires a more
387 extensive study of the northern and southern PMM dynamics that is beyond the scope of this
388 study.

389 **6. Summary and Discussion**

390 In this study, we have examined the interannual variability in the number of mesoscale
391 eddies in the SCS and its relationship with ENSO. While previous studies have already provided
392 useful findings on the impact of ENSO on SCS eddies, our research uncovered a change in the

393 ENSO influence in recent years, specifically around 2004. Furthermore, we find that the ENSO-
394 eddy relationship is controlled by two wind stress curl mechanisms: one directly related to ENSO
395 and the other related to precursors of ENSO – the northern and southern PMMs.

396 These two mechanisms appear as the two leading EOF modes of the interannual variability
397 in surface wind stress curl over the SCS. While the direct ENSO mechanism produces a negative
398 ENSO-eddy correlation throughout the analysis period, the indirect PMM mechanism produces a
399 negative ENSO-eddy correlation before 2004 but a positive ENSO-eddy correlation afterward.
400 As a result, ENSO can strongly impact the number of SCS eddies through the additive effects of
401 both the ENSO and ENSO-precursor (i.e., the PMMs) mechanisms before 2004 but has little
402 impact on the SCS eddy number after 2004 due to the cancellation between these two
403 mechanisms. The differing ENSO-eddy correlation produced by the PMM mechanism may be
404 related to a phase change of the IPO around 2004, which links the SCS eddies to the northern
405 PMM during the pre-2004 period but to the southern PMM during the post-2004 period.

406 One major finding of this study is that both the northern and southern PMMs can produce
407 strong impacts on mesoscale oceanic eddies in the SCS, which has not been documented
408 previously. While the northern PMM has been extensively studied, the southern PMM has not
409 received the same attention. Much is still unknown about why and how the relative importance
410 of these two PMMs can be modulated by the phase change of the IPO and what that implies for
411 decadal changes in SCS eddy activities. These issues were not addressed in this study and require
412 further investigations. It should be noted that findings reported here are based on datasets that are
413 rather short for a study of interannual and (especially) interdecadal variability due to the limited
414 availability of altimetry observations in the SCS. We suggest that the analyses presented here
415 should be repeated when longer data become available to verify our findings.

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644 **Figure Captions**

645 **FIG. 1.** The mesoscale eddies identified by the GV method using the SLA (contour; in cm) and
646 current (vectors; in cm/s) observed on January 26, 1993. The red and blue circles denote the
647 edges of anticyclonic and cyclonic eddies, respectively.

648 **FIG. 2.** (a) Monthly values of the number of mesoscale eddies in the SCS (solid-blue) and
649 Niño3.4 index (dashed-red) and (b) their 10-year sliding correlations. Also shown are the
650 Niño3.4 index after an 11-month running means has been applied (solid-red) in (a) and the 95%
651 and 99% significance levels of a Student-t test in (b).

652 **FIG. 3.** The spatial patterns of the EOF1 (a) and EOF2 (b) modes of the wind stress curl
653 anomalies (color and contour) over the SCS. Also shown are the regressions of surface wind
654 anomalies (vector) onto PC1 (a) and PC2 (b).

655 **FIG. 4.** The principal components (PC1 and PC2) of the first two EOF modes of wind stress curl
656 anomalies over the SCS (a) and their 10-year sliding correlations with the Niño3.4 index (b). The
657 PC values shown are normalized by their standard deviations and have been applied with an 11-
658 month running mean. The 90% and 95% significance levels are marked in (b).

659 **FIG. 5.** Lead-lagged regressions of SST (color; in °C) and sea surface wind anomalies (vectors;
660 in m/s) onto PC1 during P1 (a-d) and P2 (e-h), and onto PC2 during P1 (i-l), and P2 (m-p). The
661 lag values (in months) are shown at the left top of each panel. The white dots and red vectors
662 indicate regressions that exceed the 90% significance level, based on a Student's test.

663 **FIG. 6.** Same as Fig. 5, except for the regressions of SLP (color and shading; in Pa) and sea
664 surface wind (vector; in m/s) anomalies.

665

666 **FIG. 7.** SLP (color and contour; in Pa) and surface wind (vector; in m/s) anomalies in the SCS
667 regressed onto the PC1 index (upper) and PC2 index (lower) during sub-periods P1 (left) and P2
668 (right). The white dots and red vectors indicate regressions that exceed the 90% significance
669 level, based on a Student's test.

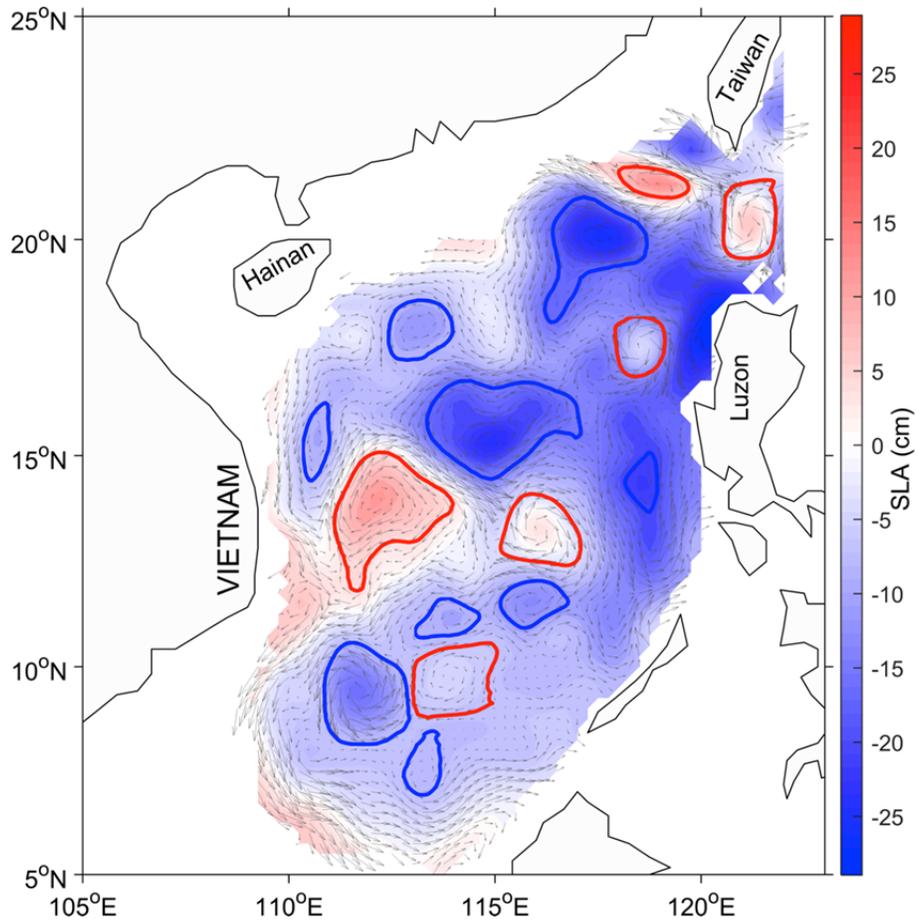
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671 northern and southern PMM indices during sub-periods P1 (left) and P2 (right). (c and d) same
672 as (a and b), except for PC2. The 90% and 95% significance levels of Student-t test are marked.

673 **FIG. 9.** SLP (color; in Pa) and surface wind anomalies (vectors; in m/s) regressed onto the
674 nPMM index (upper) and the sPMM index (sign-reversed, see text, lower) during sub-periods P1
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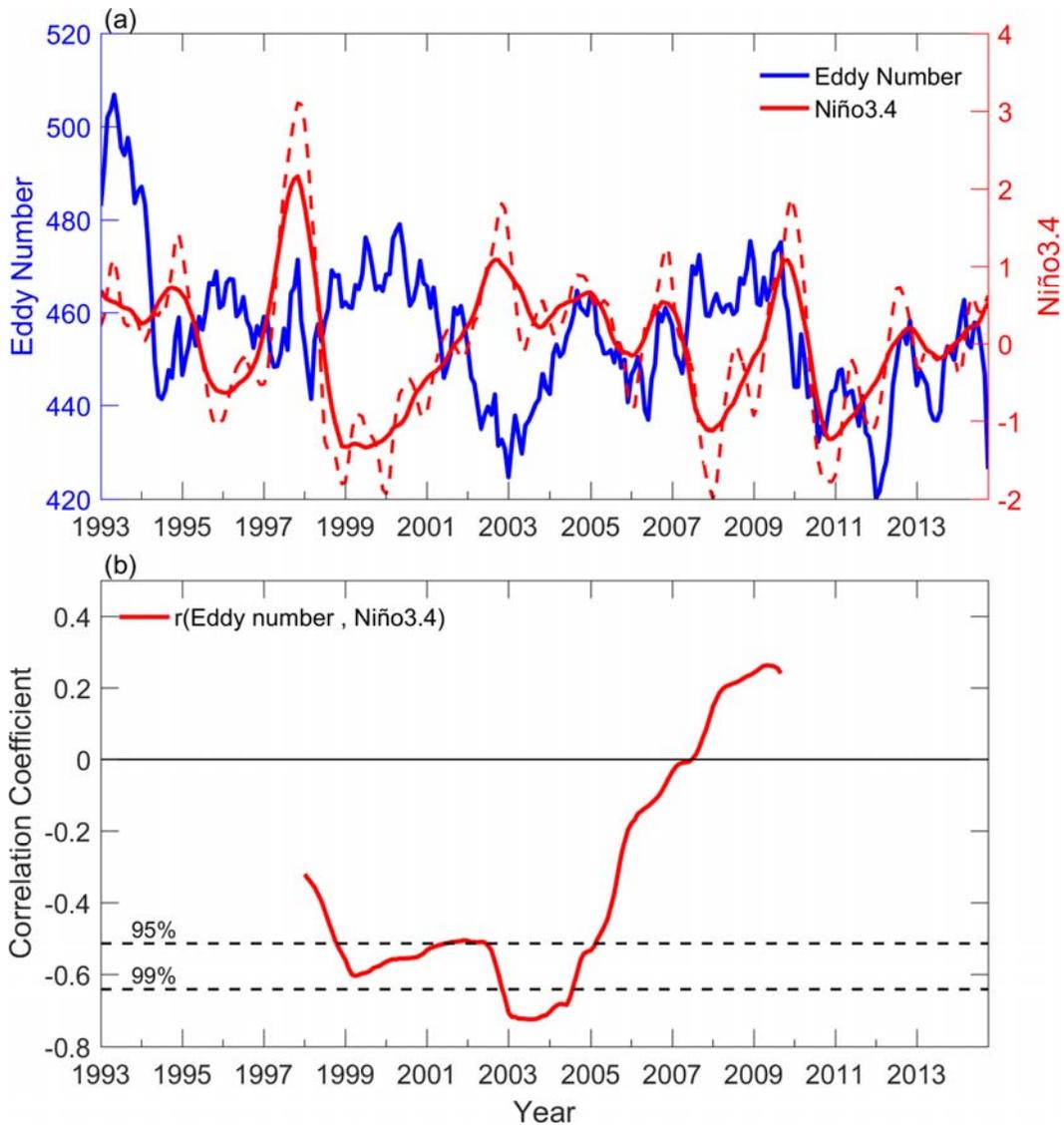
677 **FIG. 10.** (a) The 10-year sliding correlations between the eddy number anomaly in the SCS and
678 the nPMM. (b) same as (a), except for the sPMM. Also shown are the 90% and 95% significance
679 levels of a Student-t test in (a) and (b).

680 **FIG. 11.** The IPO index after a 10-year running mean in unit of standard deviation. The red-
681 colored part marks the analysis period of this study.

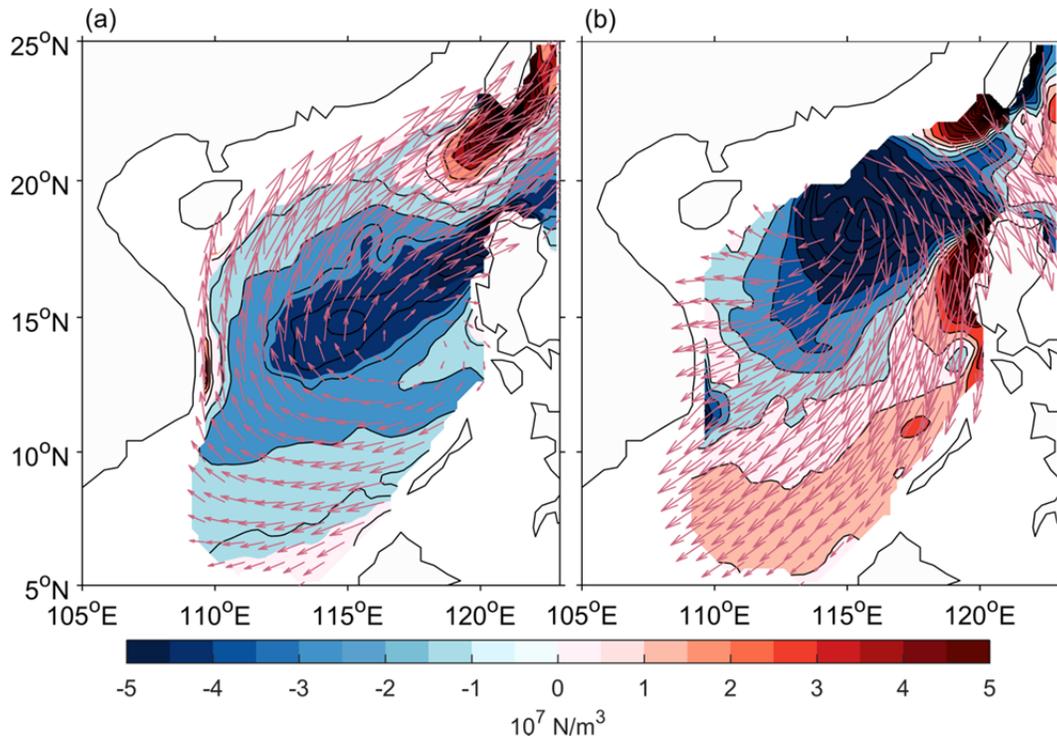
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683 during the P1 (upper) and P2 (lower) sub-periods. The regressions of surface wind (vector; in
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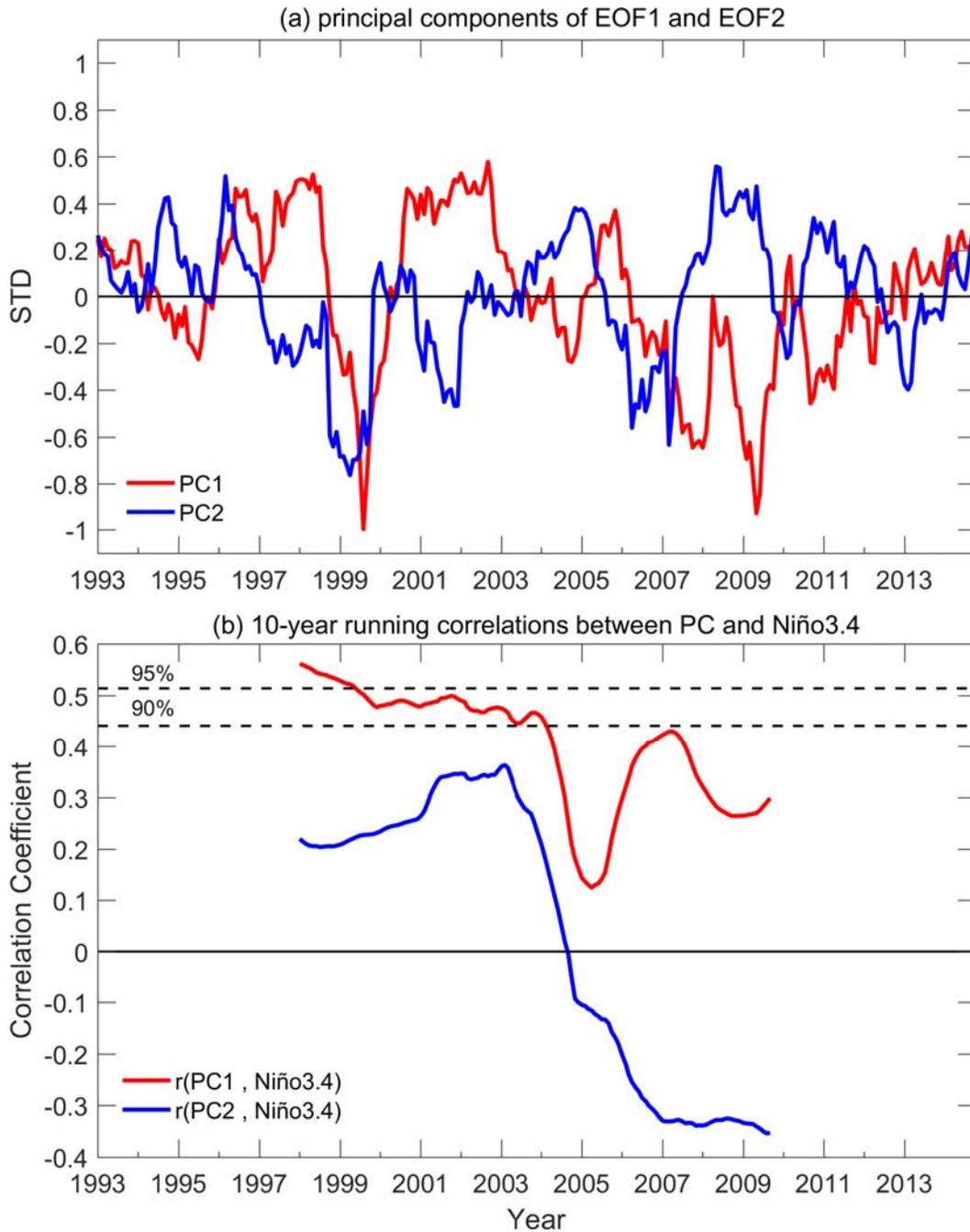
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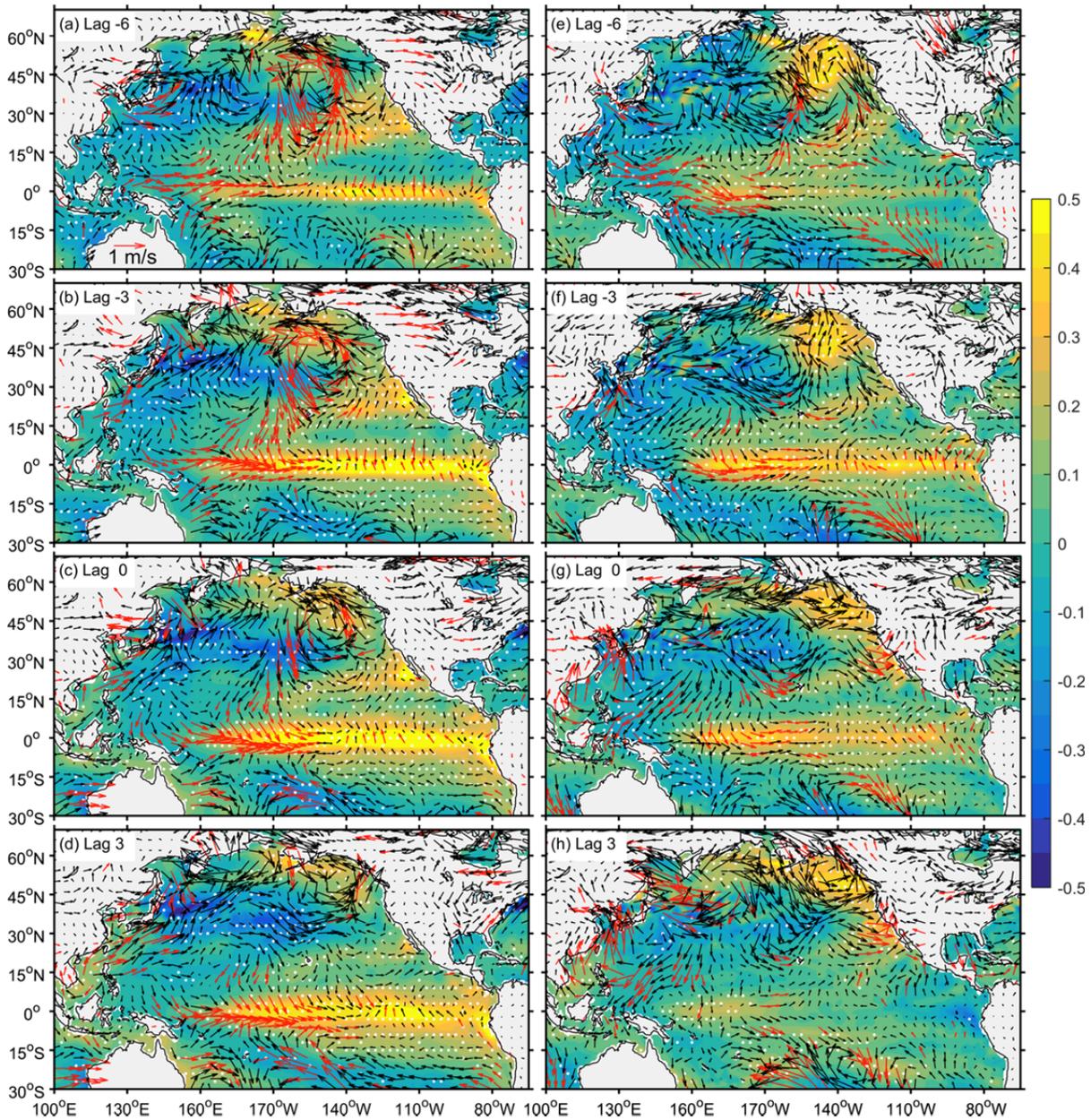
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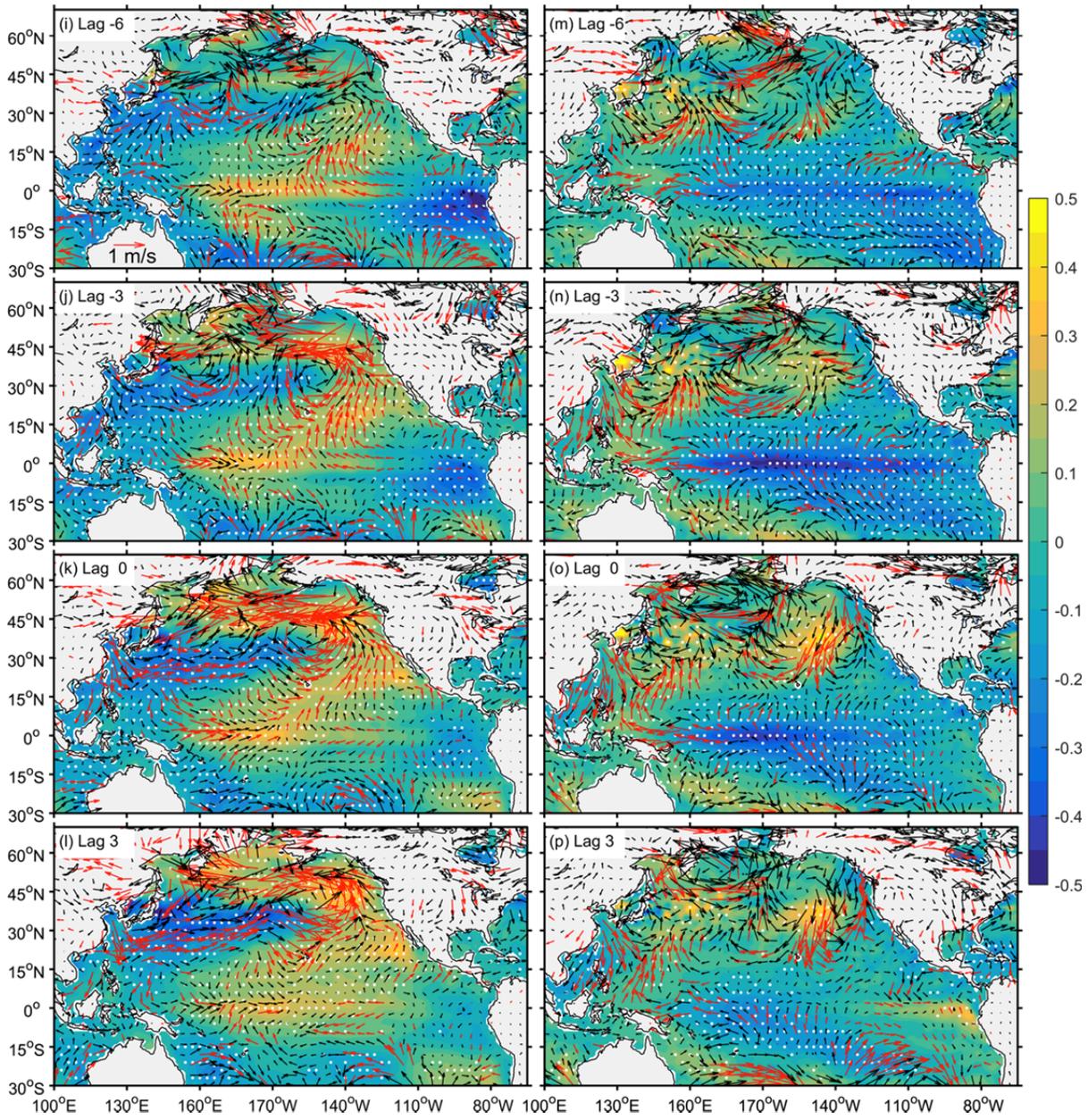
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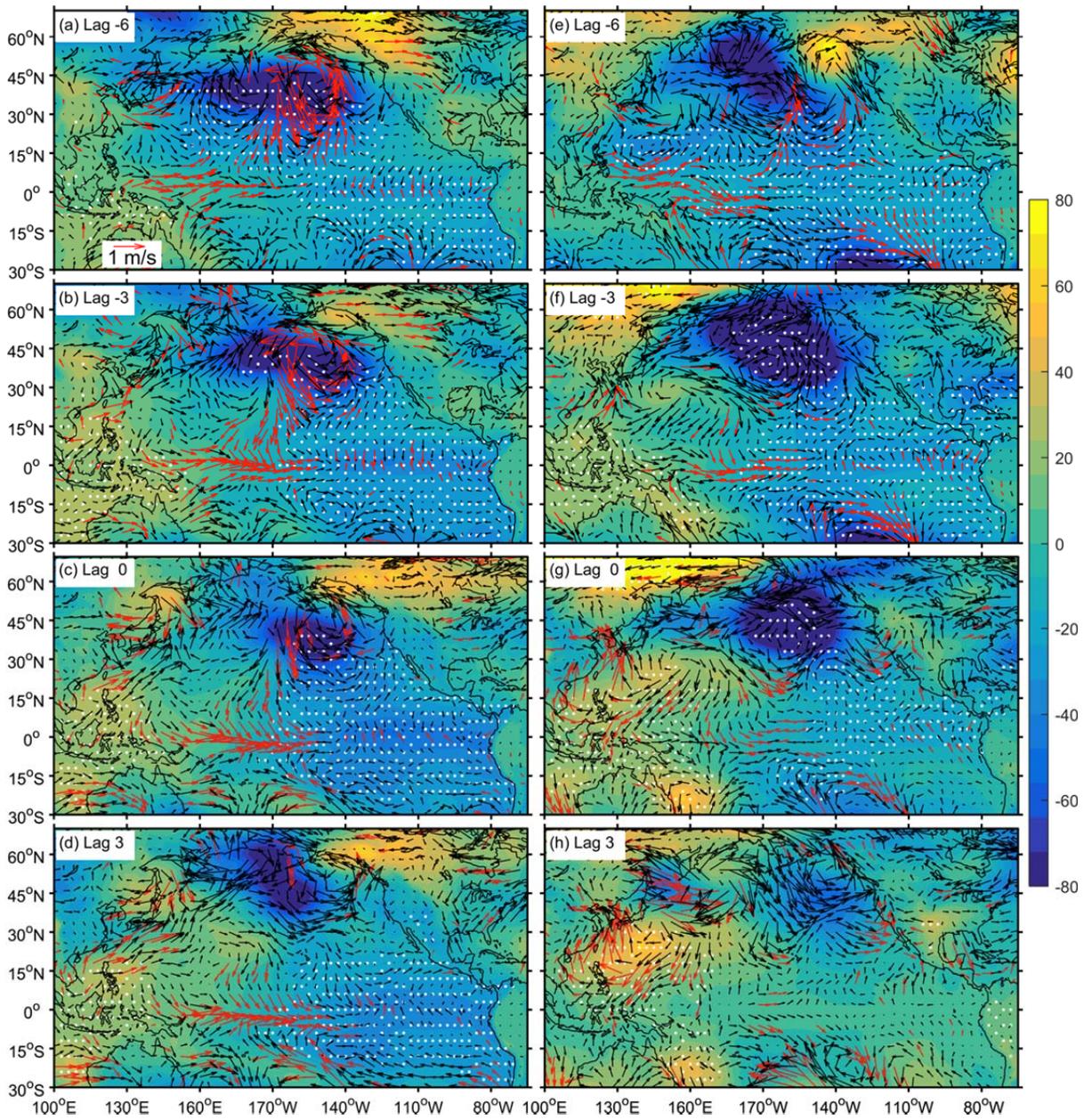
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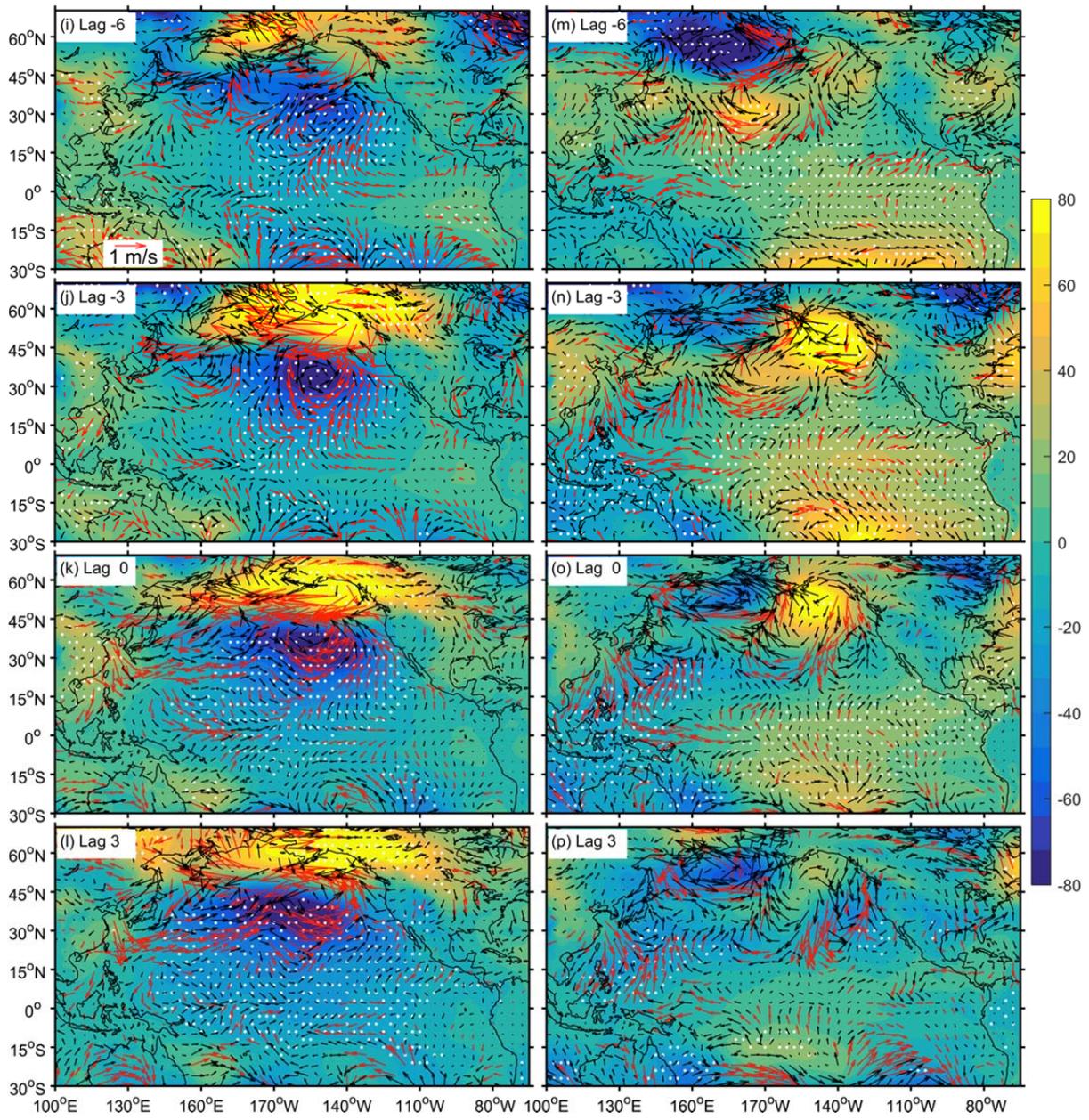
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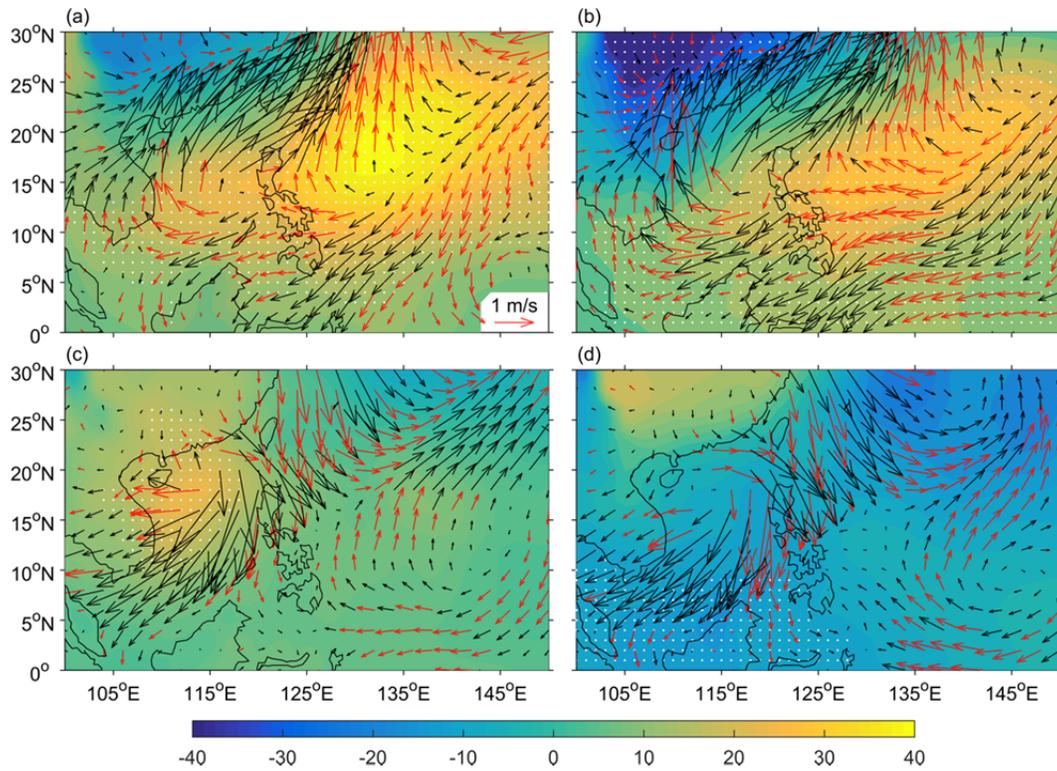
704 FIG. 5. Continued.



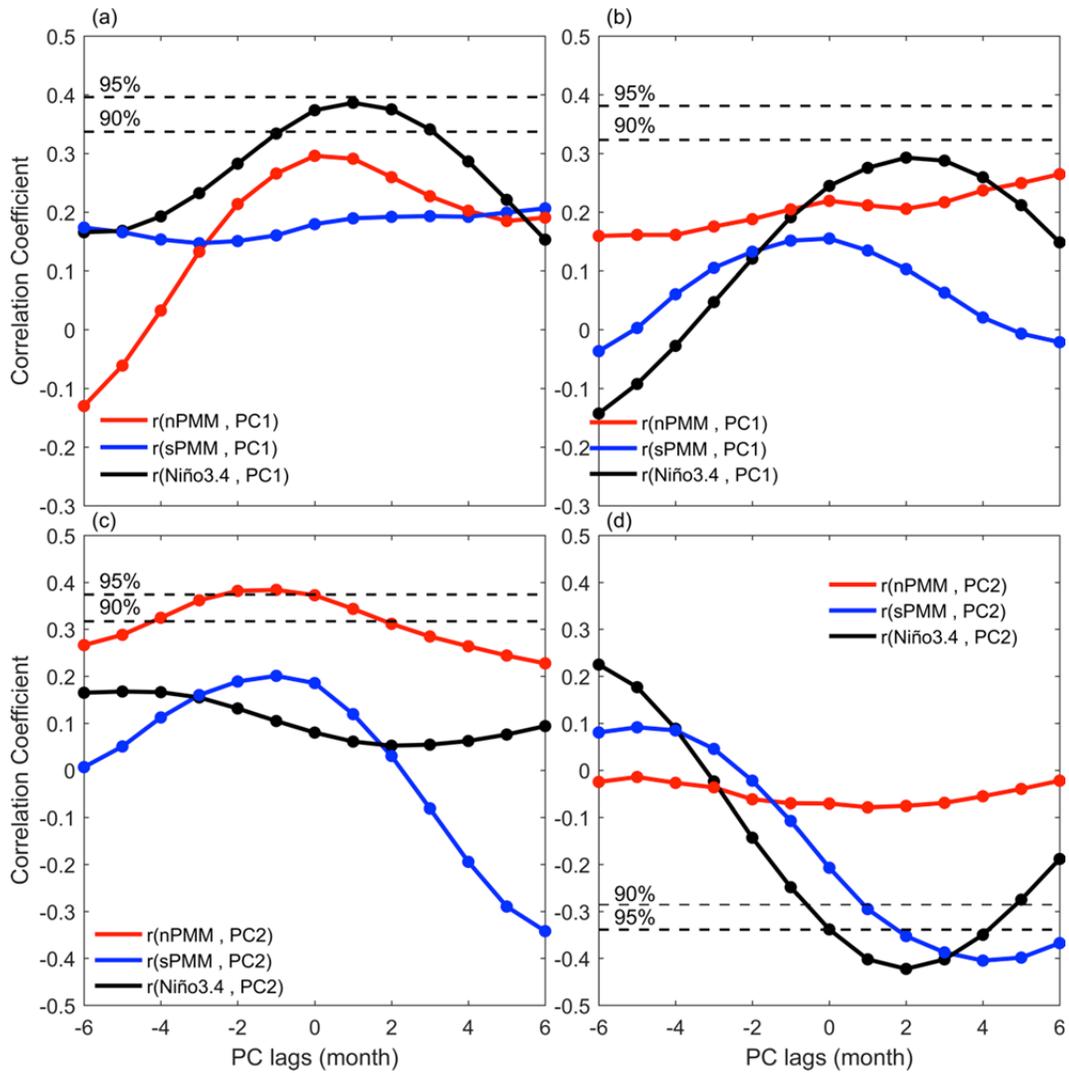
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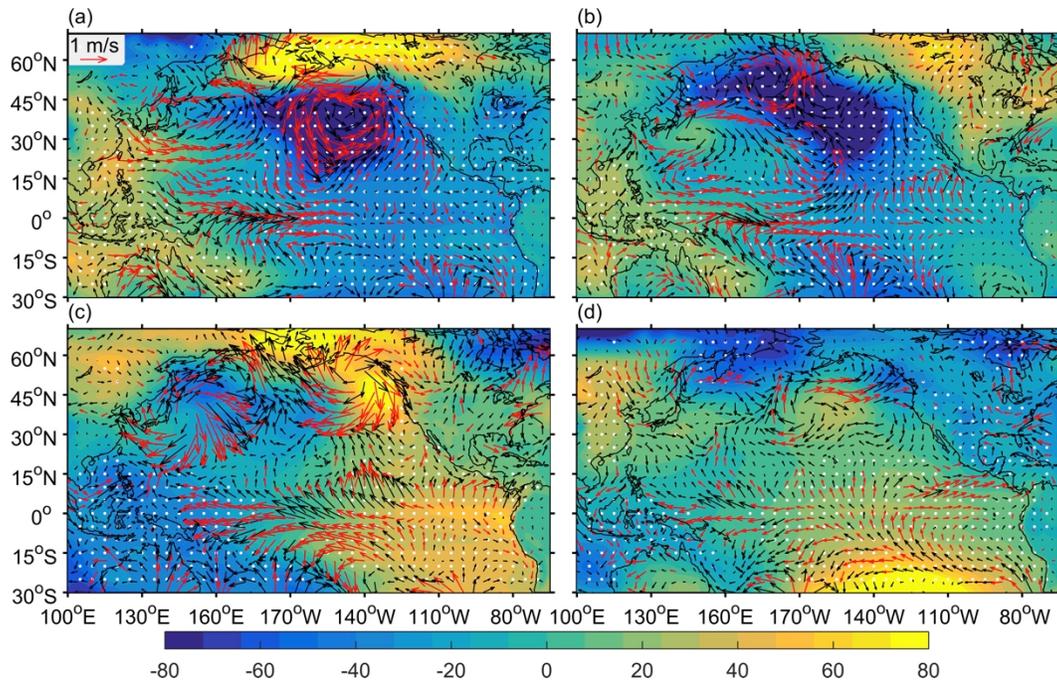
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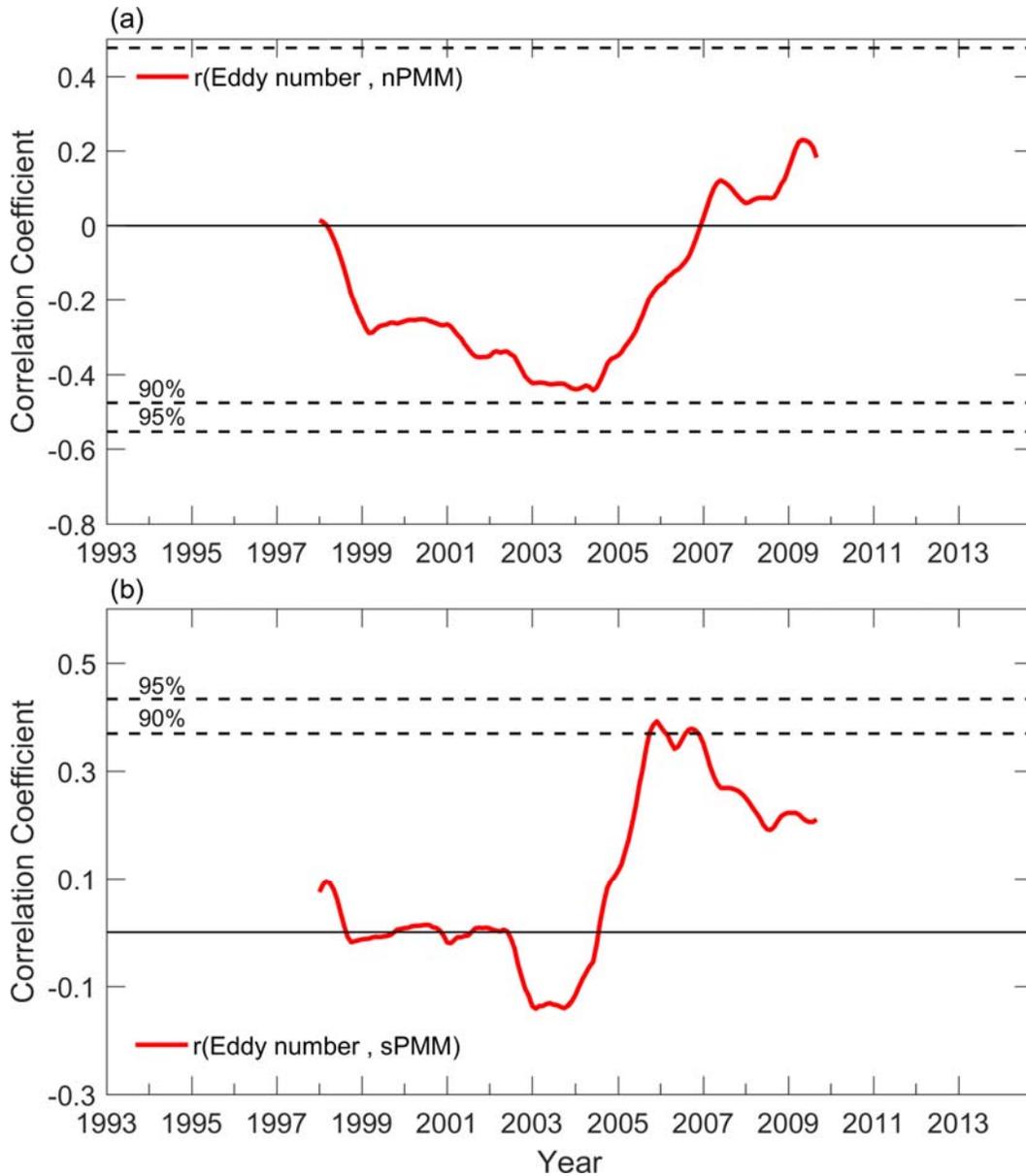
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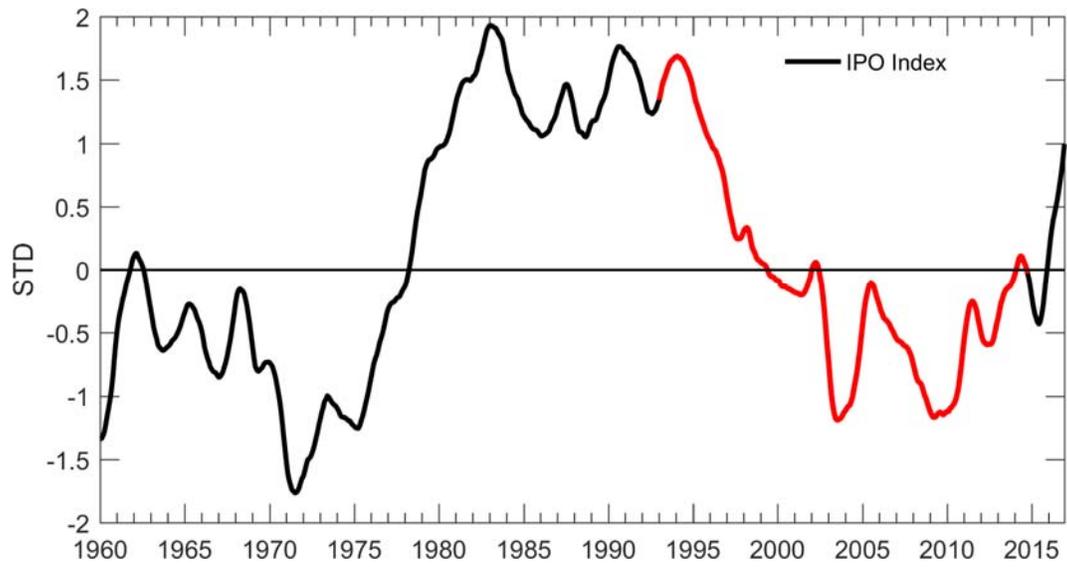
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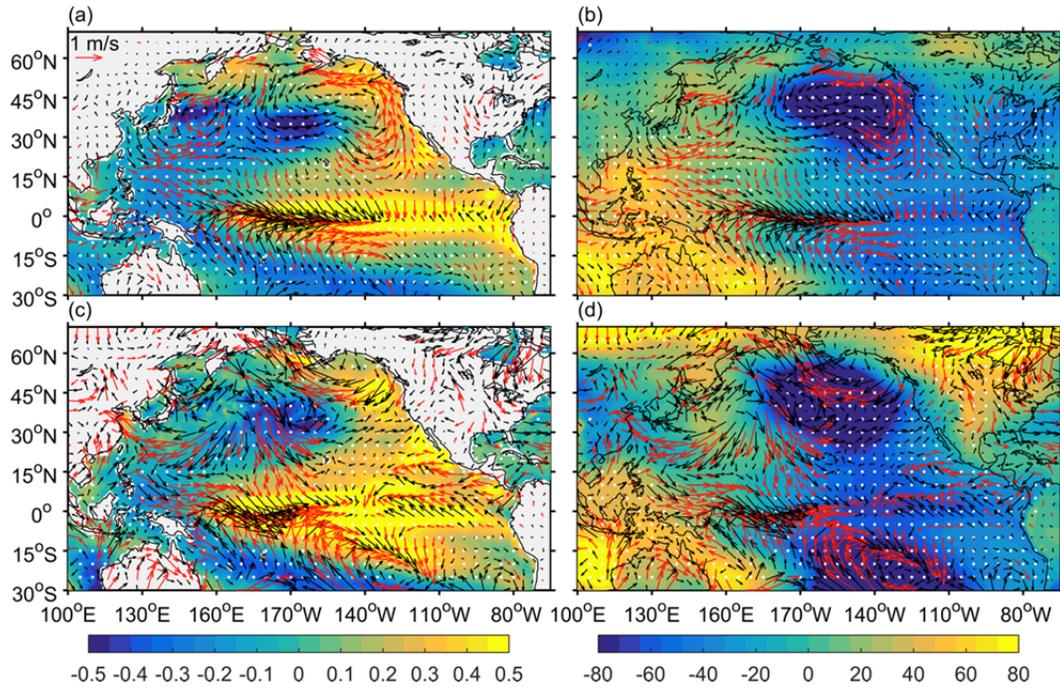
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